

# The impact of the detected gravitational waves on fundamental physics

Gerhard Schäfer

Theoretisch-Physikalisches Institut  
Friedrich-Schiller-Universität Jena

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## Outline

- The detected GWs
- BH spacetimes
- Binary dynamics
- Impact on fundamental physics

## The detected GWs

GW150914 from BBH (36/29) $M_{\odot}$  @420 Mpc; LIGO SNR 24, FAR  $6/10^7$  yrs

LVT151012 from BBH (23/13) $M_{\odot}$  @1000 Mpc; LIGO SNR 10, 37/10<sup>2</sup> yrs (1.7  $\sigma$ )

GW151226 from BBH (14/08) $M_{\odot}$  @440 Mpc; LIGO SNR 13,  $6/10^7$  yrs (5.3  $\sigma$ )

GW170104 from BBH (31/19) $M_{\odot}$  @880 Mpc (0.18); LIGO SNR 13, 1/70,000 yrs

GW170814 from BBH (31/25) $M_{\odot}$  @540 Mpc; LIGO/Virgo SNR 18, 1/27,000 yrs

GW170817 from BNS (1.5/1.3)\* $M_{\odot}$  @40 Mpc; LIGO/Virgo SNR 32, 1/80,000 yrs

GRB 170817A detected by Fermi-GBM 1.7 s after GW170817 near NGC 4993\*\*

optical SSS17a/AT 2017gfo: less than 11 hours after merger with two week long

Kilonova

\*(1.36-1.60/1.17-1.36) Low-spin priors, (1.36-2.26/0.86-1.36) High-spin priors

Chirp mass 1.188  $M_{\odot}$

\*\*  $z = 0.008$ : Hubble constant  $H_0 = 70$  (67.90 from Planck) km/sMpc

tidal deformability (significant above 600 Hz):

$$\Lambda = \frac{2}{3} k_2 \left( \frac{c^2 R}{G M} \right)^5$$

stellar radius:  $R$

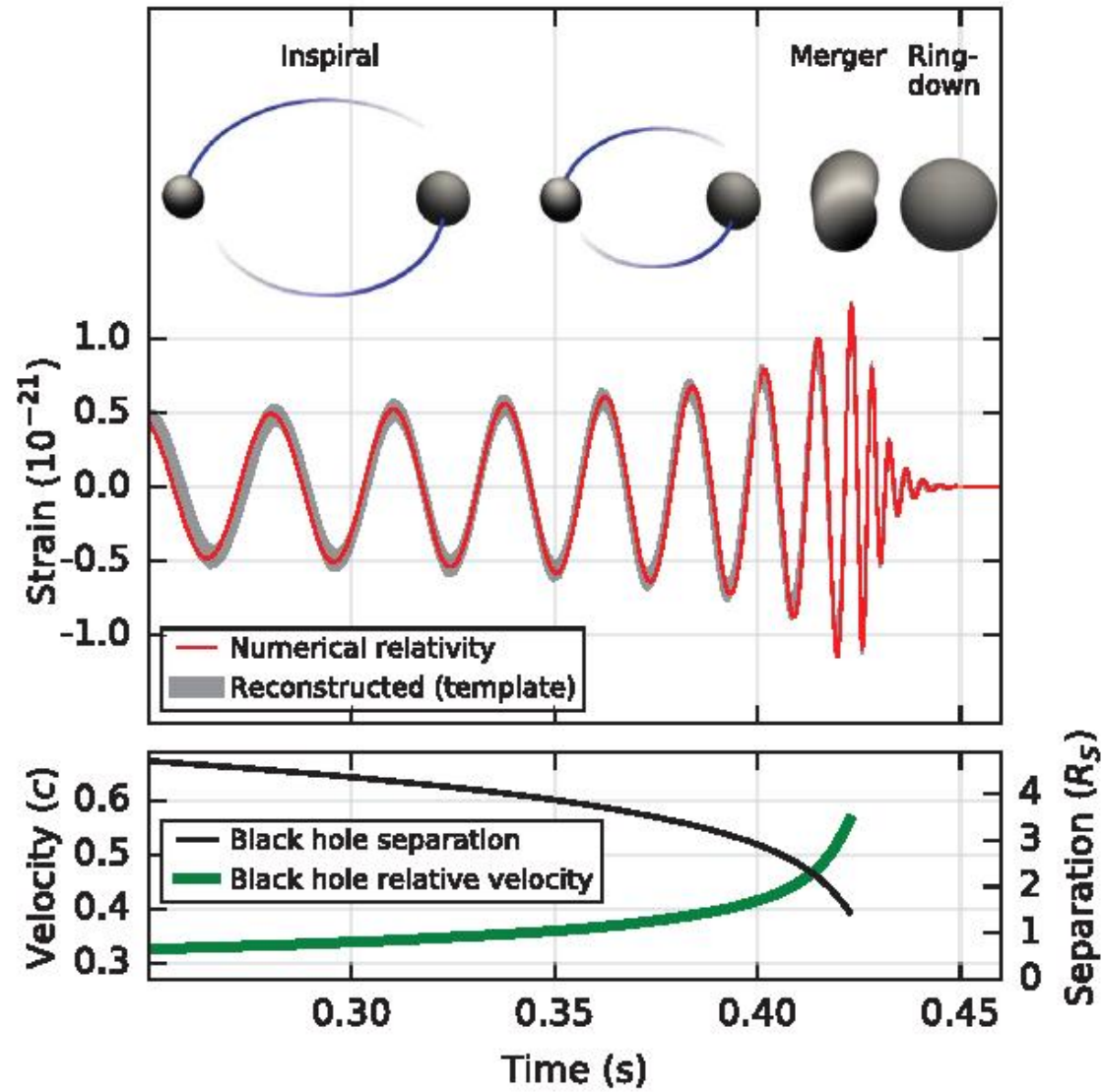
second Love number:

$$k_2 = 0.05 - 0.15 \text{ (NS)}$$

$$k_2 = 0 \text{ (BH)}$$

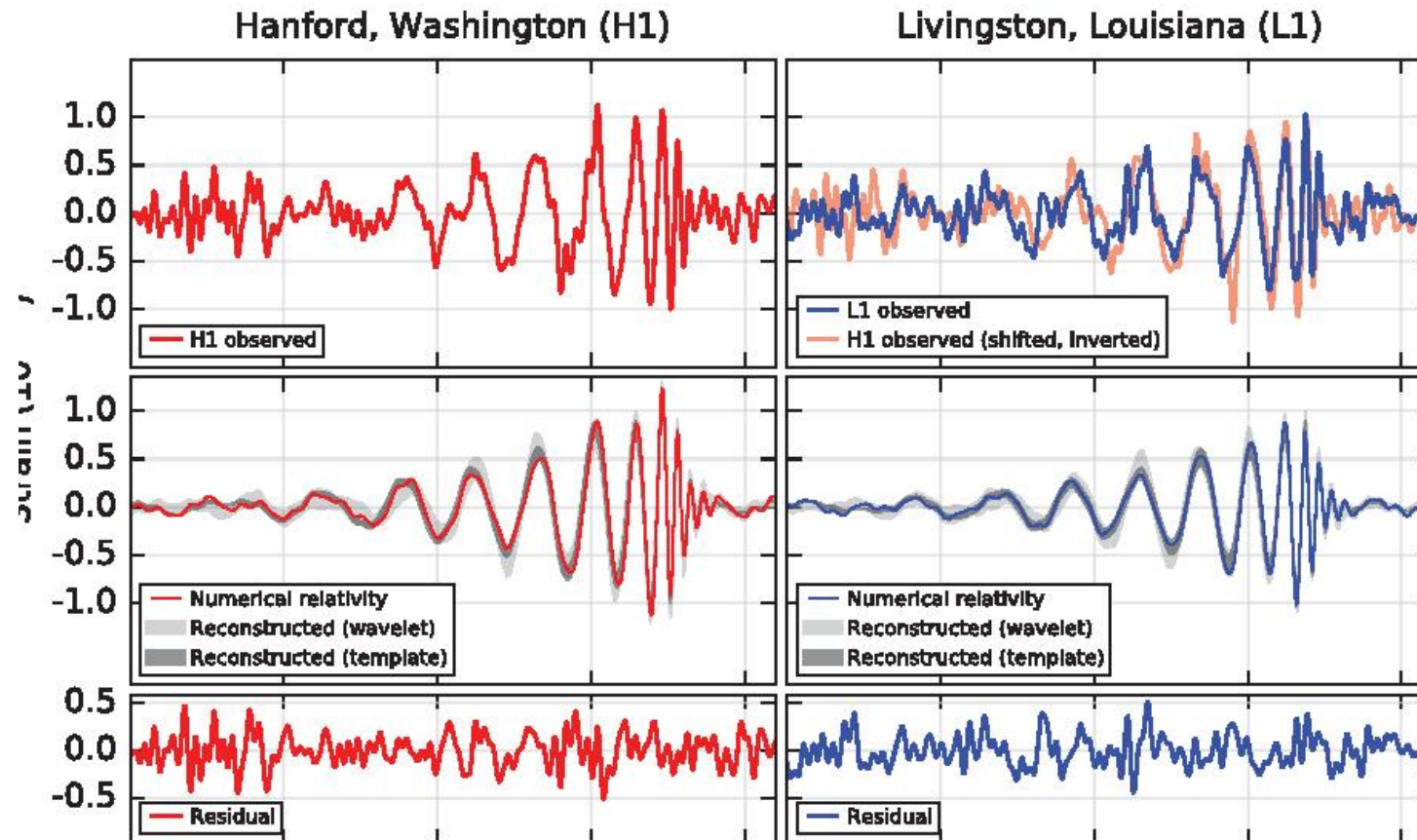
$\Lambda(1.4M_{\odot}) \leq 800$  for Low-spin priors and  $\leq 1,400$  for High-spin priors

# GW150914



$$h(t)L = \Delta L(t) = \delta L_x - \delta L_y$$

# GW150914



source power (3 solar masses radiated away in 0.2 seconds)

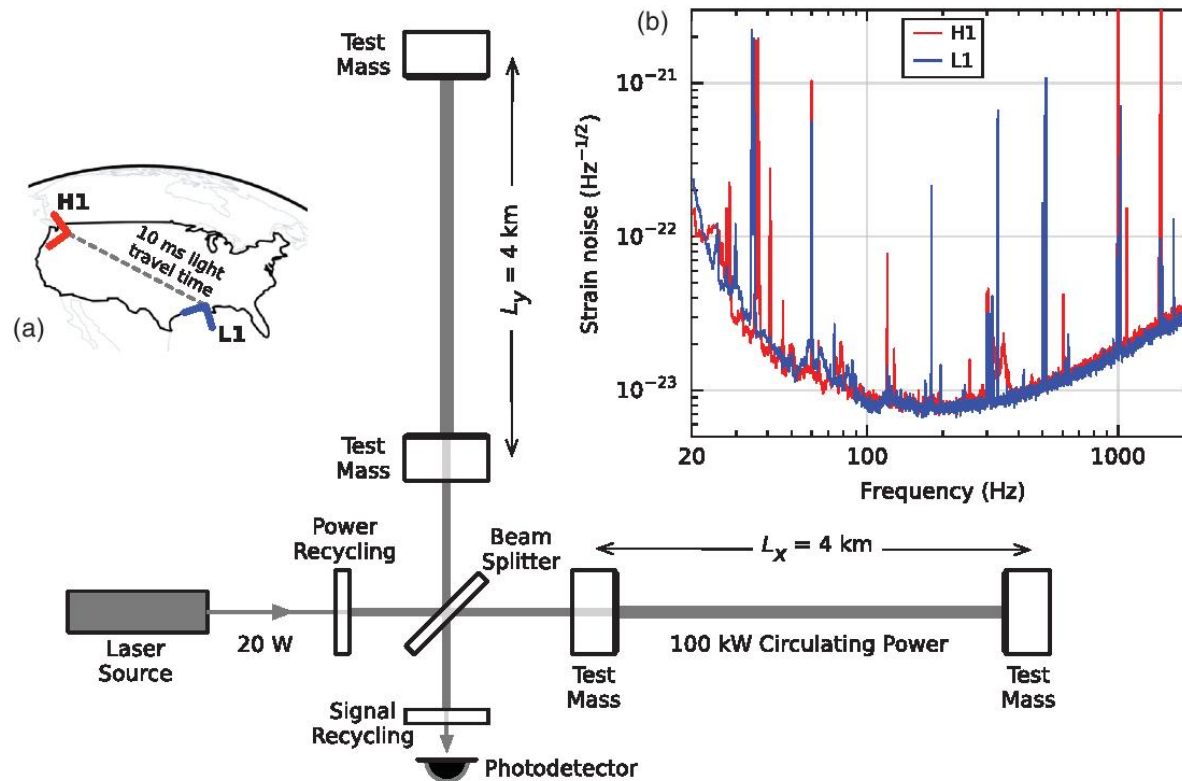
(3/1/2/2.7/0.025)

$$(6 \times 10^{30} \text{kg}) \cdot 9 \times 10^{16} \left( \frac{\text{m}}{\text{sec}} \right)^2 / 200 \text{ms} = 2.7 \times 10^{48} \text{Watt} = 3 \times 10^{-4} \frac{c^5}{4G}$$

maximum power of a single process (Dyson luminosity)

$$\frac{c^5}{4G} = \frac{Mc^2}{(4GM/c^2)/c} = 0.9 \times 10^{52} \text{Watt}$$

# Interferometric detection of GWs

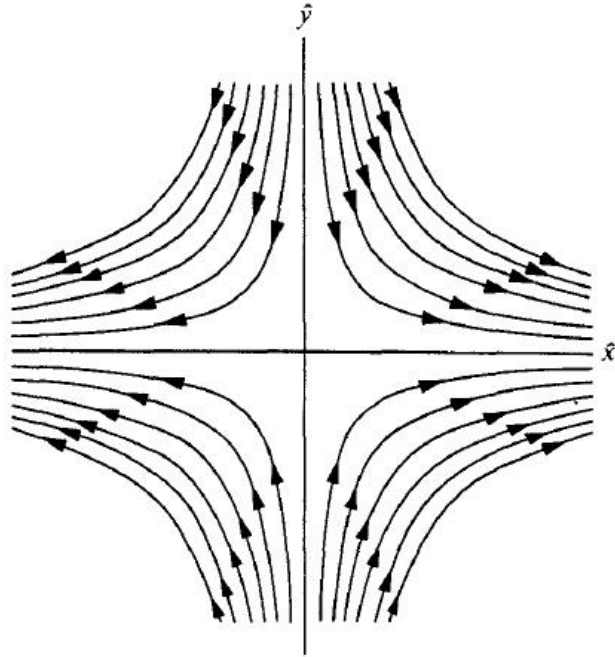




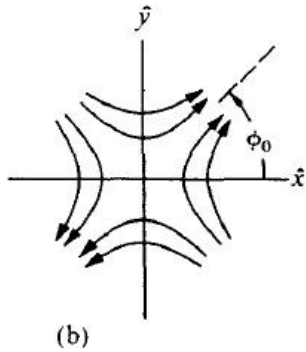
line element	$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$
gravitational wave	$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu + h_{ij}^{\text{TT}} dx^i dx^j$

Fermi normal coordinates

$$ds^2 = -\left(1 - \frac{1}{2c^2} \ddot{h}_{ij}^{\text{TT}} \hat{x}^i \hat{x}^j\right) c^2 dt^2 + d\hat{x}^i d\hat{x}^i \quad (|\hat{\mathbf{x}}| \ll c/f)$$



(a) Force lines for  $\ddot{A}_x = 0, \ddot{A}_+ > 0$



(b)

gravitational quadrupole wave

propagating in  $\hat{z}$ -direction

$$d^2 \hat{x} / dt^2 = \frac{1}{2} (\ddot{A}_+ \hat{x} + \ddot{A}_\times \hat{y})$$

$$d^2 \hat{y} / dt^2 = \frac{1}{2} (-\ddot{A}_+ \hat{y} + \ddot{A}_\times \hat{x})$$

$$d^2 \hat{z} / dt^2 = 0$$

MTW: Gravitation (1973)

chirp mass

$$\mathcal{M} := \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} = \frac{c^3}{G} \left( \frac{5}{96} \frac{\dot{f}}{\pi^{8/3} f^{11/3}} \right)^{3/5}$$

velocity

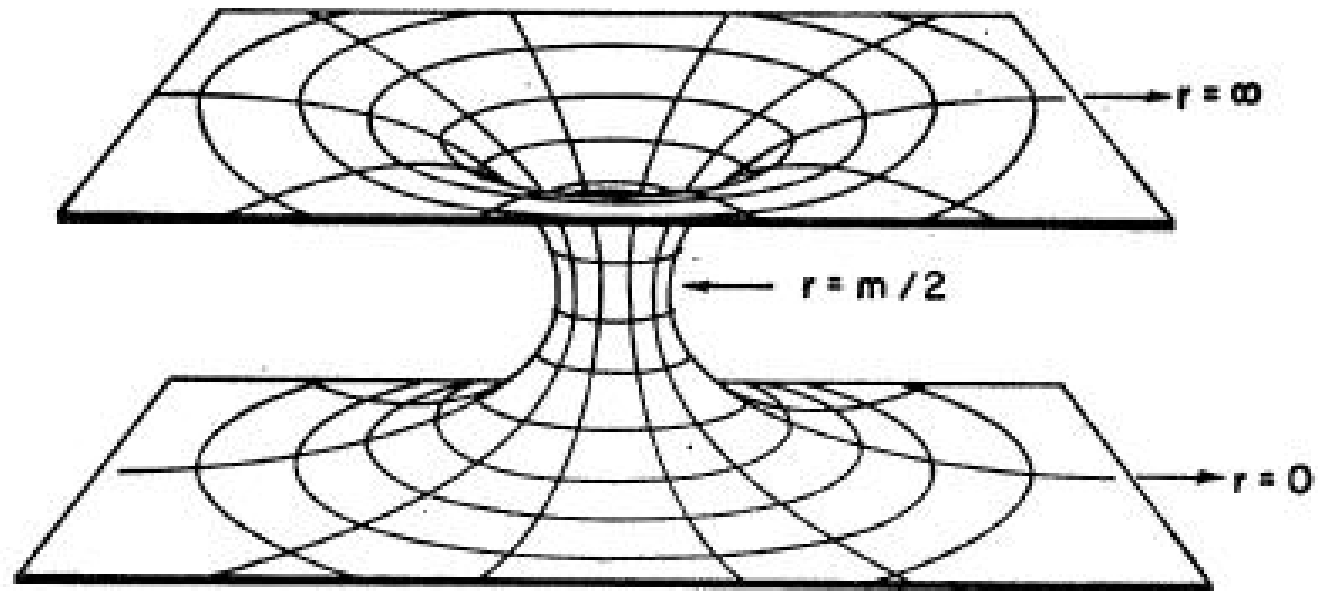
$$\frac{v}{c} = \left( \frac{GM\pi f}{c^3} \right)^{1/3}$$

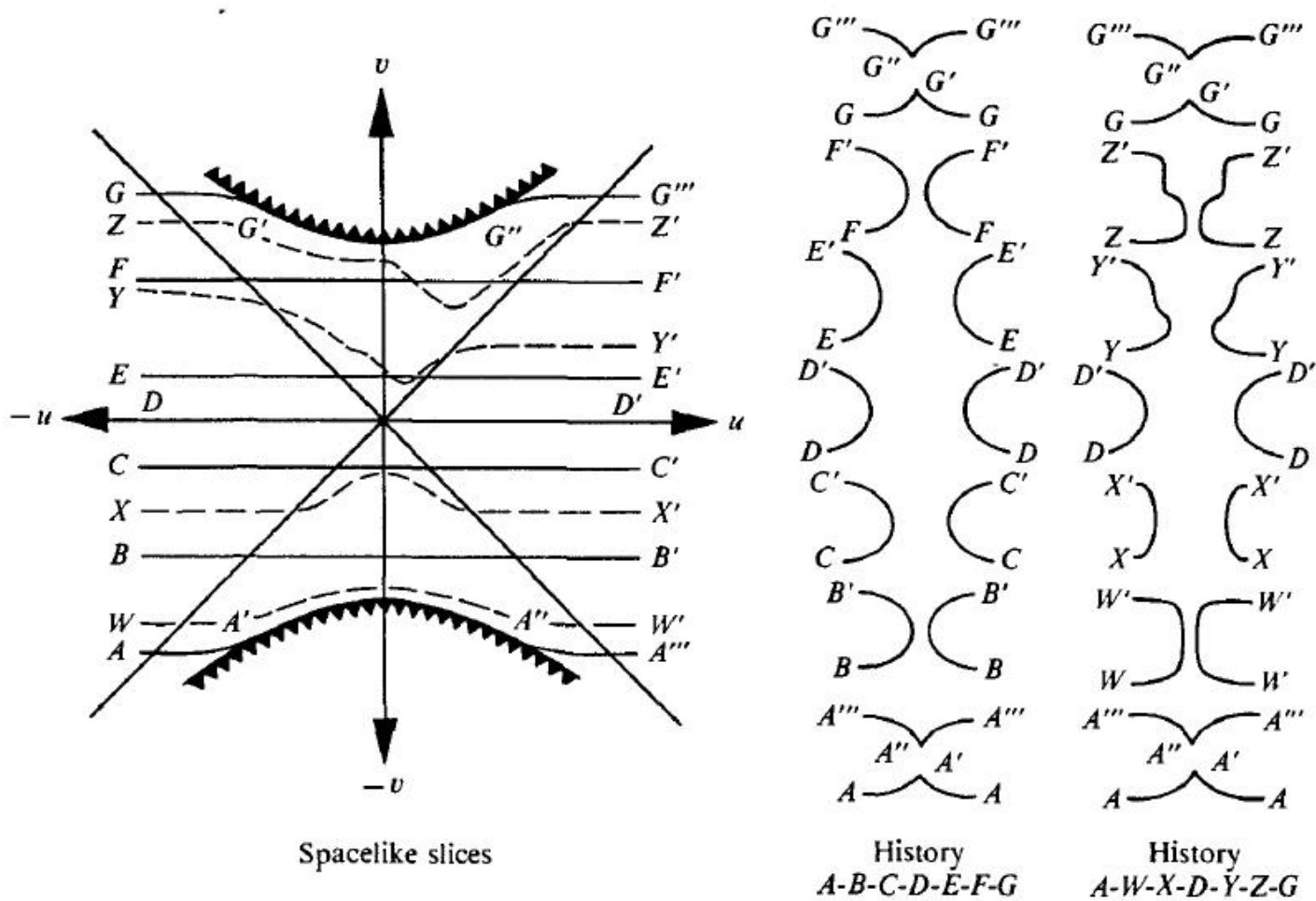
gravitational quadrupole wave

$$h_{ij}^{\text{TT}}(t, \mathbf{x}) = \frac{2}{r} \frac{G}{c^4} P_{ijkl}(\mathbf{n}) \ddot{Q}_{km} \left( t - \frac{r}{c} \right)$$

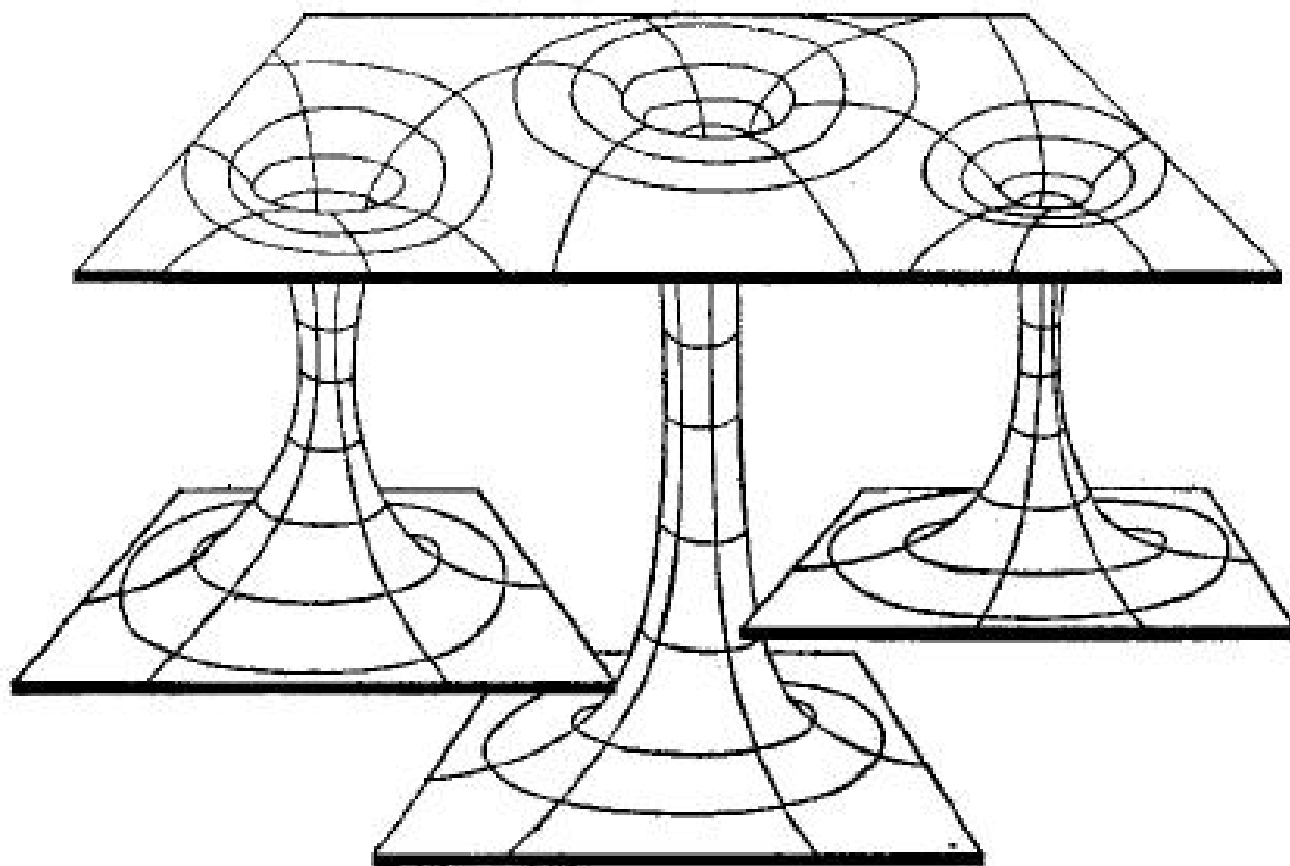
$$Q_{km} = \sum_a m_a (x_k^a x_m^a - \frac{1}{3} \delta_{km} x_l^a x_l^a)$$

# BH spacetimes (non-rotating)

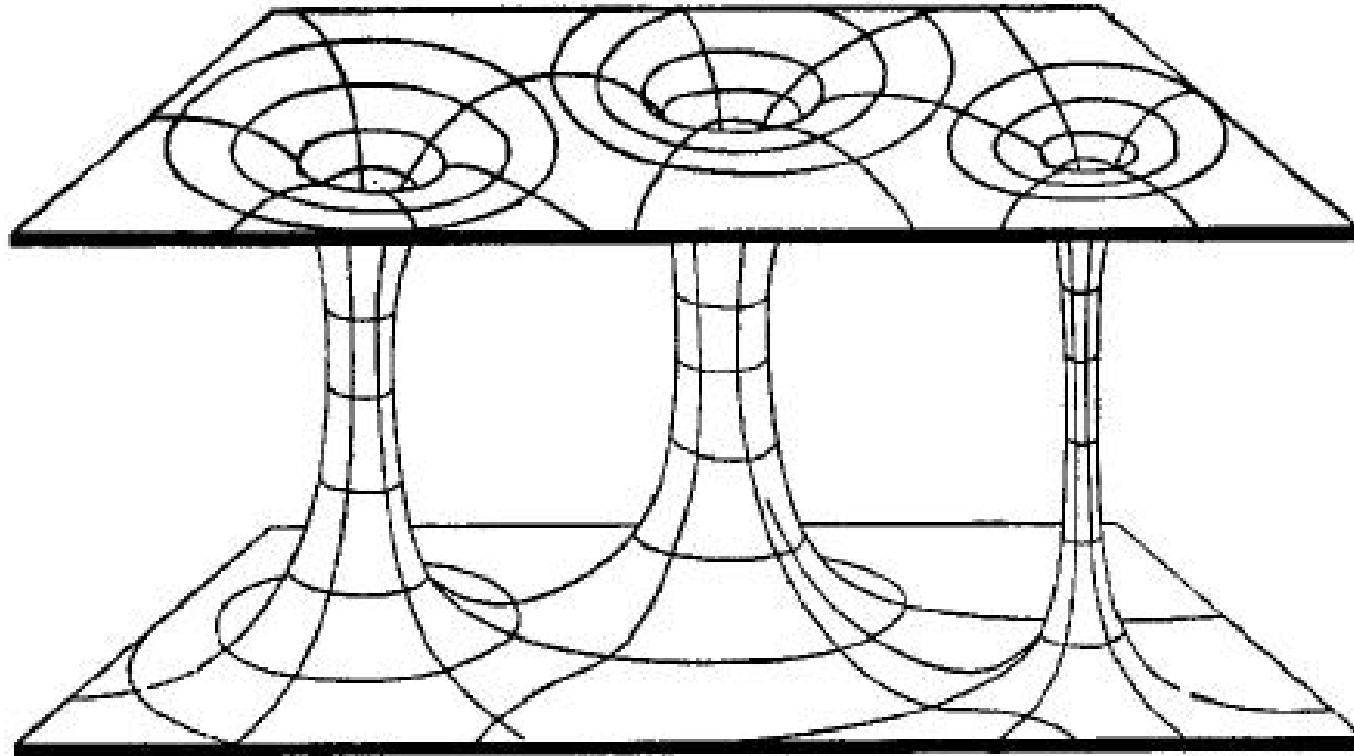




MTW: Gravitation (1973)



Brill-Lindquist BHs



Misner-Lindquist BHs

## Binary dynamics

Test particle in Schwarzschild field

$$r = 6 GM/c^2 = 3 R_S \quad (x = 1/6): \text{ ISCO}$$

$$r = 3 GM/c^2 = 3/2 R_S \quad (x = 1/3): \text{ light ring or light sphere}$$

$$r = 2 GM/c^2 = R_S \quad (x = 1/2): \text{ event horizon}$$

$$r = 2.25 GM/c^2 = 9/8 R_S: \text{ Buchdahl limit (most compact star)}$$



## Innermost Stable Circular Orbit (ISCO)

$$H = H(\mathbf{p}, \mathbf{r}), \quad p^2 = p_r^2 + j^2/r^2, \quad p_r = (\mathbf{p} \cdot \mathbf{r})/r$$

$$\text{circular orbits: } p_r = 0, \quad p^2 = j^2/r^2, \quad H = H(j, r)$$

$$\text{circular motion: } \frac{\partial}{\partial r} H(j, r) = 0 \rightarrow H(j)$$

$$\text{orbital frequency: } \omega = \frac{dH(j)}{dj} \rightarrow H(\omega)$$

$$\text{ISCO: } \frac{\partial^2}{\partial r^2} H(j, r) = 0 \text{ or alternatively } \frac{dH(\omega)}{d\omega} = 0$$

$$E(x) \equiv \frac{H(x)}{\mu c^2}, \quad x = \frac{(GM\omega)^{2/3}}{c^2}$$

$$\begin{aligned} \text{SBH: } E(x) &= \frac{1-2x}{\sqrt{1-3x}} = \sqrt{1-x \frac{1-4x}{1-3x}} \\ &= 1 - \frac{1}{2}x + \frac{3}{8}x^2 + \frac{27}{16}x^3 + \frac{675}{128}x^4 + \frac{3969}{256}x^5 + \dots \end{aligned}$$

## Binaries at 4PN approximation:

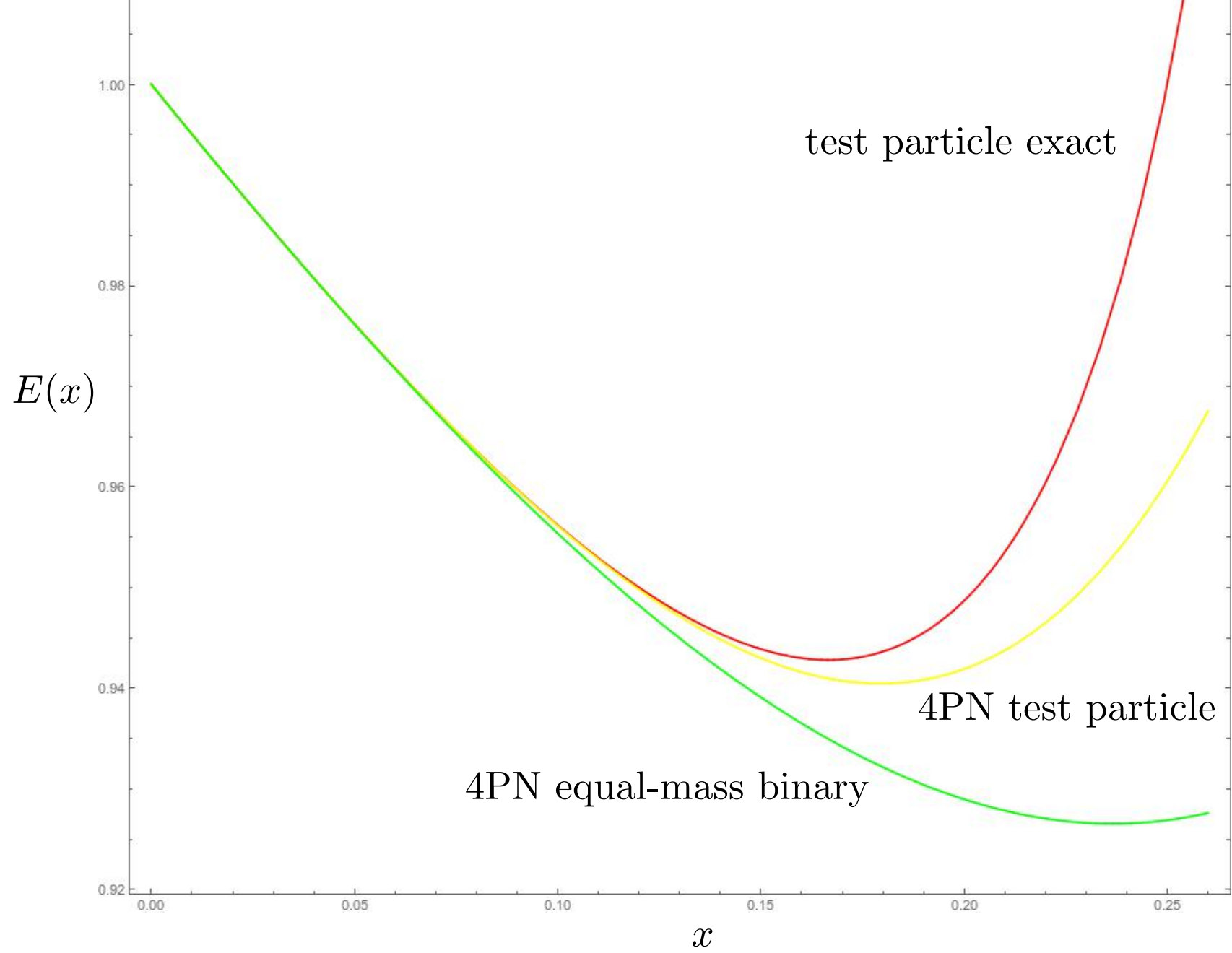
Jaranowski/GS 2013, Damour/Jaranowski/GS 2014 (see PRD or arXiv)

$$\begin{aligned} E_{4PN}(x) &= 1 - \frac{x}{2} + \left( \frac{3}{8} + \frac{1}{24} \nu \right) x^2 + \left( \frac{27}{16} - \frac{19}{16} \nu + \frac{1}{48} \nu^2 \right) x^3 \\ &+ \left( \frac{675}{128} + \left( -\frac{34445}{1152} + \frac{205}{192} \pi^2 \right) \nu + \frac{155}{192} \nu^2 + \frac{35}{10368} \nu^3 \right) x^4 \\ &- \frac{1}{2} \left( -\frac{3960}{128} + (c_1 + \frac{448}{15} \ln x) \nu + \left( -\frac{498449}{3456} + \frac{3157}{576} \pi^2 \right) \nu^2 + \frac{301}{1728} \nu^3 + \frac{77}{31104} \nu^4 \right) x^5 \end{aligned}$$

$$c_1 = -\frac{123671}{5760} + \frac{9037}{1536} \pi^2 + \frac{1792}{15} \ln 2 + \frac{896}{15} \gamma_E$$

$$\nu = \frac{\mu}{M} = \frac{m_1 m_2}{M^2}, \quad M = m_1 + m_2, \quad \mu = \frac{m_1 m_2}{M}$$

$$0 \leq \nu \leq \frac{1}{4} = 0.25, \quad \nu = 0: \text{ test particle}, \quad \nu = 0.25: \text{ equal masses}$$



## Matched Filtering

S. Finn: *Detection, measurement, and gravitational radiation*,  
PRD 46, 5236 (1992)

C. Cutler, E. Flanagan: *Gravitational waves from merging compact binaries:  
How accurately can one extract the binary's parameters from the inspiral  
waveform?*, PRD 49, 2658 (1994)

A. Królak, K. Kokkotas, G. Schäfer: *Estimation of the post-Newtonian  
parameters in the gravitational-wave emission of a coalescing binary*,  
PRD 52, 2089 (1995)

J. Creighton, W. Anderson: *Gravitational-Wave Physics and Astronomy*,  
WILEY-VCH Verlag, Weinheim 2011

data stream:  $s(t; \theta) = n(t) + h(t; \tilde{\theta})$ ,  $\theta = (\theta_i)$ ,  $i = 1, \dots, N$

$n(t)$ : noise amplitude

$h_{te} = h(t; \hat{\theta})$ : template wave form

inner product:

$$\langle h_{te} | h \rangle = 2 \int_0^\infty \frac{\tilde{h}_{te}^*(f) \tilde{h}(f) + \tilde{h}^*(f) \tilde{h}_{te}(f)}{S_n(f)} df = 4 \operatorname{Re} \int_0^\infty \frac{\tilde{h}_{te}^*(f) \tilde{h}(f)}{S_n(f)} df$$

$S_n(f)$ : spectral noise density,  $S_n(f) = \lim_{T \rightarrow \infty} \frac{1}{T} |\tilde{n}(f; T)|^2$

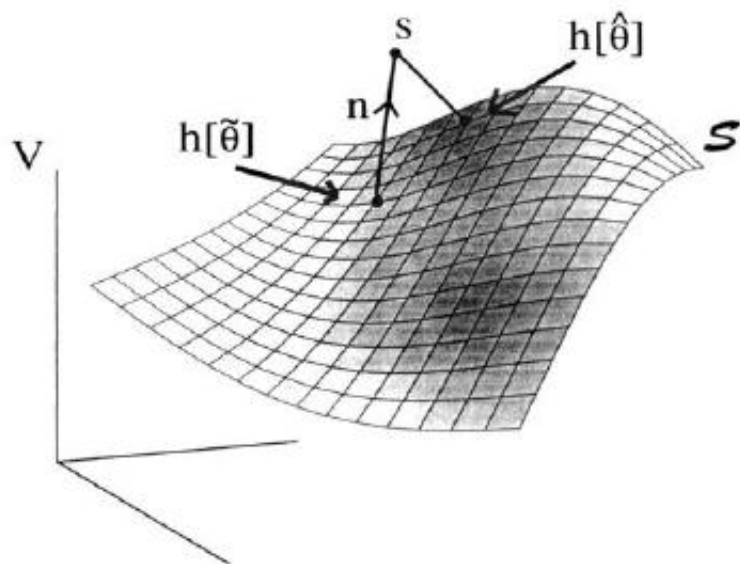


FIG. 1. Gravitational waveforms from coalescing compact binaries are completely specified by a finite number of parameters  $\theta = (\theta^1, \dots, \theta^k)$ , and so form a surface  $S$  in the vector space  $V$  of all possible measured detector outputs  $s = s(t)$ . The statistical properties of the detector noise endow  $V$  with the structure of an infinite-dimensional Euclidean space. This figure illustrates the relationships between the true gravitational wave signal  $h(\tilde{\theta})$ , the measured signal  $s$ , and the “best-fit” signal  $h(\hat{\theta})$ . Given a measured detector output  $s = h(\tilde{\theta}) + n$ , where  $n = n(t)$  is the detector noise, the most likely values  $\hat{\theta}$  of the binaries parameters are just those that correspond to the point  $h(\hat{\theta})$  on the surface  $S$  which is closest [in the Euclidean distance  $(s - h | s - h)$ ] to  $y$ .

$$\text{SNR} : \frac{S}{N}[h_{te}] = \frac{\langle h_{te}|h_{te} \rangle}{\text{rms} \langle h_{te}|n \rangle} = \langle h_{te}|h_{te} \rangle^{1/2}$$

1.5PN wave form (including tail from back scattering) in Fourier domain above ISCO ( $c = G = 1$ ):

$$\tilde{h}_{te}(f) = \frac{Q}{r} \mathcal{M}^{5/6} f^{-7/6} \exp[i\Psi(f)] \quad 0 < f < (6^{3/2}\pi M)^{-1}$$

$$\Psi(f, \hat{\theta}) =$$

$$2\pi f t_c - \phi_c - \frac{\pi}{4} + \frac{3}{4} (8\pi \mathcal{M} f)^{-5/3} \left[ 1 + \frac{20}{9} \left( \frac{743}{336} + \frac{11}{4} \nu \right) x - 16\pi x^{3/2} \right]$$

+ spin-orbit + spin1-spin2 + rotational quadrupole ( $\text{spin}^2$ ) + tidal quadrupole

## Impact on fundamental physics

- 1 - GW do exist! a la Einstein? longtime observations through BNS!
- 2 - BH do exist! a la Einstein? multipole structure? alternatives?
- 3 - NS-components likely! proper rotations? EOS? alternatives?