# Compensation of gravity gradients and rotations in precision atom interferometry



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partly based on Phys. Rev. Lett. 118, 160401 (2017)

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## Motivation





 $m_{i}$ 

- Central to Einstein's equivalence principle.
- Tests of UFF with macroscopic masses:
  - ▶ free fall, lunar laser ranging (LLR)

• torsion balance (Ëotvös)  
= 
$$m_{\rm g}$$

Why atoms?

$$\eta_{AB} = 2 \frac{|g_A - g_B|}{g_A + g_B} \lesssim 10^{-13} \dots 10^{-15}$$

#### Measurement of Newton's gravitational constant G

• By far the less accurately determined of all fundamental constants (using *macroscopic* masses).



#### Earth observations: mapping Earth's gravitational field

- Gravity gradiometry from space (spatial resoultion  $\geq 100 \, \mathrm{km}$ )
- Observing time evolution of mass distribution with applications to geophysics, hydrology, oceanography ...

Example: ground water depletion and stressed aquifers



Atom interferometry can make a significant contribution in all these cases

### BUT

*Gravity gradients* (and *rotations*) pose a major challenge due to the dependence on the *initial position* and *velocity* of the atomic wave packets.

### Outline

- I. Motivation
- 2. Precise gravitational measurements with atom interferometry:
  - atom-interferometry-based gravimeters
  - long-time interferometry (& microgravity platforms)
- 3. Challenges in UFF tests due to gravity gradients
- 4. Overcoming loss of contrast and initial co-location problem
- 5. Compensation of large *rotation* rates

# Precise gravitational measurements with atom interferometry (AI)

#### Al-based gravimeters



 $\mathbf{k}_{\mathrm{eff}} = \mathbf{k}_1 - \mathbf{k}_2$ 

$$\delta \phi = - k_{\rm eff} \ g \, T^2$$

 $\int_{0}^{\frac{1}{2}} \int_{\pi} \int_{2\pi} \int_{3\pi} \delta\phi$ 

 $N_{{\rm g}_1}/(N_{{\rm g}_1}+N_{{\rm g}_2})$ 

The evolution of the wave packets can be decomposed into two independent aspects:

expansion dynamics of a centered wave packet

central position and momentum which follow classical trajectories including the kicks from the laser pulses

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#### Gradiometry and measurements of G



Fixler et al., Science **315**, 74 (2007)

- differential measurement
- common-mode noise suppression
- determination of the gravity gradient

$$\Gamma_{zz} = -\frac{\partial^2 U}{\partial z \, \partial z} \approx -\frac{g_2 - g_1}{z_2 - z_1}$$

- changing position of well-characterized source mass  $\longrightarrow$  measurement of G

 $\Delta G/G \approx 1.5 \times 10^{-4}$ 

Rosi et al., Nature **510**, 518 (2014)

Higher sensitivity --> long-time interferometry

 $\delta \phi = k_{\rm eff} \, a \, T^2$ 

- Natural compact set-ups in microgravity platforms (freely falling frame)
- Challenges:
  - ▶ growing size of atom cloud → BECs, atomic lensing
  - rotations
  - gravity gradients (effects grow cubically with time)

#### Microgravity platforms

 $\delta g \sim 10^{-5} g - 10^{-6} g$ 

drop tower in Bremen (> 500 drops) sounding rocket (23 Jan 2017) International Space Station (late 2017-)







QUANTUS (5-10s)

MAIUS (6 min)

CAL / BECCAL (several years)

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Challenges in UFF tests due to gravity gradients

#### Initial co-location

• Systematics associated with initial central position & momentum of the two species can a *mimic violation* of UFF:



No limitation in principle, but challenging in practice.
Minimum time for verification set by Heisenberg's uncertainty principle.

Minimum time for verification set by Heisenberg's uncertainty princip  $n N \sigma_p \sigma_z \ge \hbar/2$  (time required may exceed

entire mission lifetime)

#### Loss of contrast

• Gravity gradients (*tidal forces*) lead to open interferometers:



Relative displacement between interfering wave packets at exit port
fringe pattern in density profile -> loss of contrast

 $\delta z \neq 0, \ \delta p \neq 0$ 



# Overcoming loss of contrast and initial co-location problem

• Phase shift contribution connected with initial co-location directly related to  $\delta z$ ,  $\delta p$ :

$$\delta\phi = \frac{1}{\hbar}\,\delta\mathbf{p}\cdot\left(\mathbf{x}_0 + 2\,\mathbf{v}_0T\right) - \frac{1}{\hbar}\,\delta\mathbf{x}\cdot m\mathbf{v}_0 + \dots$$

• Suitable adjustment of laser wavelength of 2nd pulse  $\longrightarrow \delta z = \delta p = 0$ 



A. Roura, *Phys. Rev. Lett.* **118**, 160401 (2017)

Initial co-location as well as loss of contrast

are simultaneously taken care of.

- Required single-photon frequency change:
  - for longer times in space

 $T = 5 \,\mathrm{s} \longrightarrow \Delta \nu \sim 14 \,\mathrm{GHz}$ 

• for moderate times (and higher  $k_{eff} = 2n k_{ph}$ ) AOMs may be sufficient

2n = 50  $T = 1 \,\mathrm{s} \longrightarrow \Delta \nu \sim 0.6 \,\mathrm{GHz}$ 

• Dependence on the mirror position, but highly suppressed in the differential measurement.

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#### Circumventing Heisenberg's Uncertainty Principle in Atom Interferometry Tests of the Equivalence Principle

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- Besides tests of UFF, application to gradiometry measurements: (relaxing coupling of static  $\Gamma$  to initial-position/velocity jitter & bias)
  - mapping of Earth's gravitational field from space
  - accurate measurements of G
  - gravitational antennas

- Several other groups have also expressed great interest: Stanford, Florence, SYRTE and LKB/ENS (Paris)
- Atomic fountain experiments in Stanford's 10-meter tower



- gravity-gradient compensation scheme successfully implemented
- very effective in overcoming the initial-colocation problem
- key ingredient in efforts to test UFF with atom interferometry at 10<sup>-14</sup> level

#### Gradiometry & determination of G

- One can use the technique to cancel the effect of static gravity gradients in measurement of time-dependent ones.
- Also for measurements of static gravity gradients insensitive to initial position & velocity:



vanishing gradiometry phase for

$$\Delta \nu = \frac{c}{4\pi} \left( \Gamma_{zz} T^2 / 2 \right) k_{\text{eff}}$$

G. D'Amico et al. (submitted)

(application to determination of G)

Compensation of large rotation rates • Compensation of rotations with a tip-tilt mirror as seen from a *non-rotating frame*:



• Tip-tilt mirror leads to change in  $k_{eff}$  along the *longitudinal* direction:

 $k_{\rm eff} \rightarrow k'_{\rm eff} = \cos(\Omega T) k_{\rm eff}$ 

• It can be compensated with following change for the second pulse:

 $\Delta k_{\rm eff} \approx -(1/2)(\Omega T)^2 k_{\rm eff}$ 



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• Simultaneous compensation of gravity gradients and large rotation rates:



$$\Delta k_{\rm eff} \,=\, \left(\Gamma_{zz} \,T^2/2\right) k_{\rm eff} \,-\, (1/2) (\Omega \,T)^2 \,k_{\rm eff}$$

- Quantitative example for BECCAL
  - ► ISS's angular velocity and gravity gradient:

 $\Omega_{\rm ISS} \approx 1.13 \,\rm mrad/s$ 

 $\Gamma \approx 2.5 \times 10^{-6} \mathrm{s}^{-2}$ 

• required frequency change for  $T = 2.6 \,\mathrm{s}$ :

 $\Delta \nu \approx 1.5 \,\mathrm{GHz}$ 

(partial cancellation)

# Conclusion

- Gradiometry and tests of UFF based on AI can provide a useful complement to those based on macroscopic masses.
- Gravity gradients pose a great challenge in practice as well as an ultimate limitation from HUP due to:
  - initial co-location of the two species
  - loss of contrast
- I have presented a novel scheme that can simultaneously overcome both difficulties.
- It can be combined with a *tip-tilt mirror* to compensate large rotation rates.



### Thank you for your attention.

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