

NEW NON-UNIFORM BLACK STRING SOLUTIONS

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- Introduction
- Black Strings in $D = 5, 6$
- Topology Changing Transition?
- Black Strings in EMD Theory
- Conclusions

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*Supported by DFG

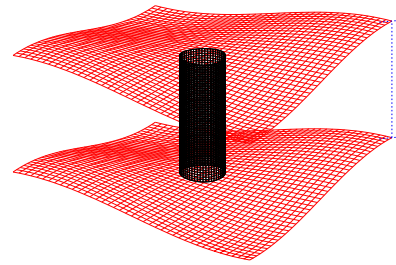
Introduction

- **vacuum solutions** in $D \geq 5$ space-time dimensions
- $(D - 1)$ space-time \times 1 extra (compact) dimension



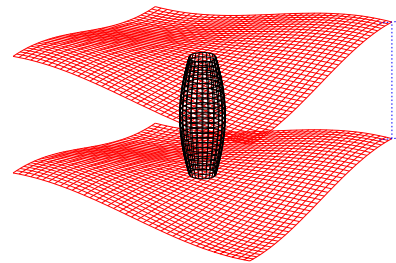
uniform black string

- $D - 1$ black hole
- translational invariance
- horizon $S^{D-3} \times S^1$
- exist for all L, M (mass)
- unstable for small M
(Gregory-Laflamme)



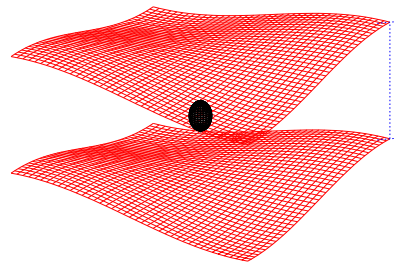
non-uniform black string

- perturbative solutions
(Gubser, Sorkin)
- depend on extra coord.
- horizon $S^{D-3} \times S^1$
- exist above GL point
- exact (num) solutions
(Wiseman, Kol)



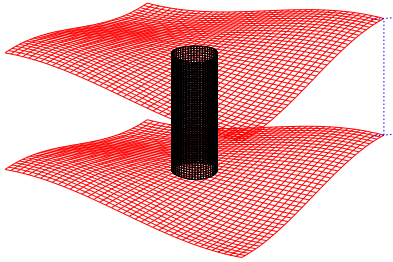
caged black hole

- horizon S^{D-2}
(Sorkin, Kol, Piran;
Kudoh, Wiseman)
- exist for small masses



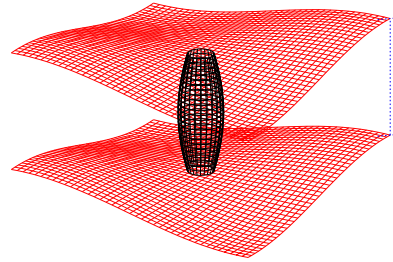
Topology changing transition?

uniform black string



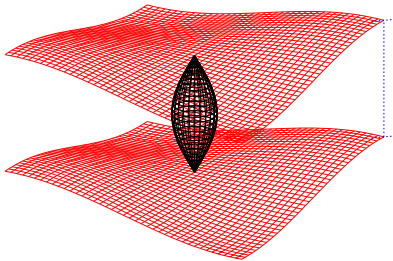
$$S^{D-3} \times S^1$$

non-uniform black string



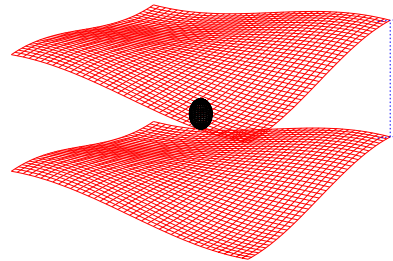
$$S^{D-3} \times S^1$$

pinch off



$$\rightarrow S^{D-2}$$

caged black hole



$$S^{D-2}$$

Ansatz

$$ds^2 = -e^{2A(r,z)} f(r) dt^2 + e^{2B(r,z)} \left[\frac{dr^2}{f(r)} + dz^2 \right] + e^{2C(r,z)} r^2 d\Omega_{D-3}^2$$

periodic coordinate z , length L

$$r^2 = \sum_{i=1}^{D-2} (x^i)^2, \quad x^D = t, \quad f(r) = 1 - \left(\frac{r_0}{r} \right)^{D-4}$$

- **Einstein equations:** $G_t^t = 0, \quad G_r^r + G_z^z = 0, \quad G_\theta^\theta = 0$
- **Constraints:** $G_r^z = 0, \quad G_r^r - G_z^z = 0$
- all other equations hold trivially
- **uniform black strings:** $A = 0, \quad B = 0, \quad C = 0$
- **Horizon at $r = r_0$**

$$\text{Transform } r = \sqrt{r_0^2 + \tilde{r}^2} \quad \rightsquigarrow \quad \text{Horizon at } \tilde{r} = 0$$

- **Boundary conditions**

$$\tilde{r} = 0 \quad A - B = d_0, \quad \partial_{\tilde{r}} A = \partial_{\tilde{r}} C = 0 \quad (\text{also } \partial_{\tilde{r}} B \text{ required})$$

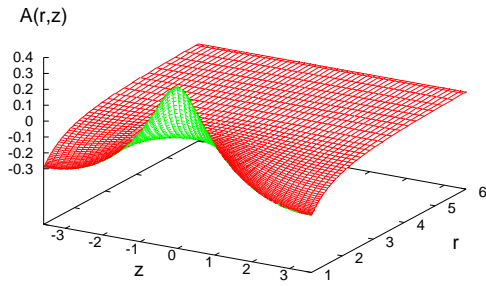
$$\tilde{r} \rightarrow \infty \quad A = B = C = 0 \quad (M^{D-1} \times S^1)$$

$$z = 0, L/2 \quad \partial_z A = \partial_z B = \partial_z C = 0$$

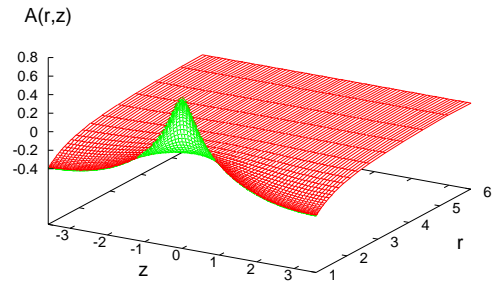
- $d_0 = \text{const.}, e^{d_0} \sim T$ **Hawking temperatur**

$$\lambda = \frac{1}{2} \left(\frac{\mathcal{R}_{\max}}{\mathcal{R}_{\min}} - 1 \right)$$

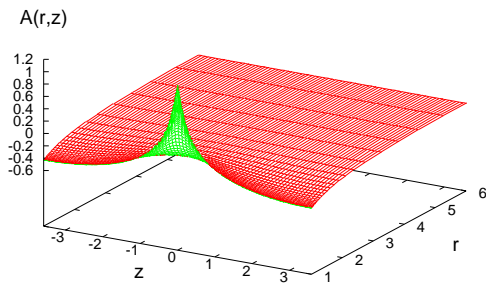
$\lambda=1.0$



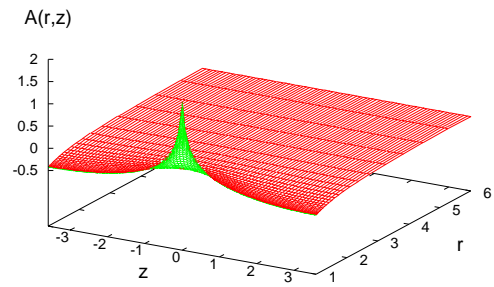
$\lambda=2.0$



$\lambda=5.0$

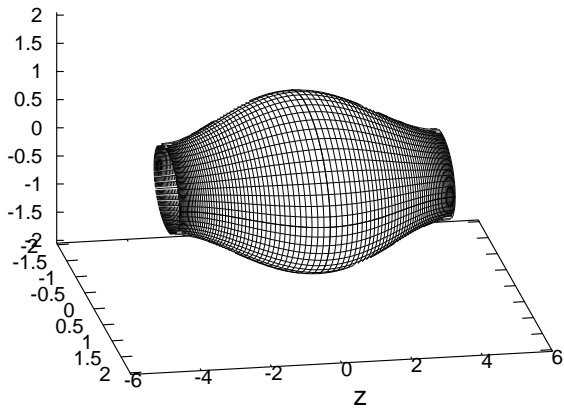


$\lambda=9.0$

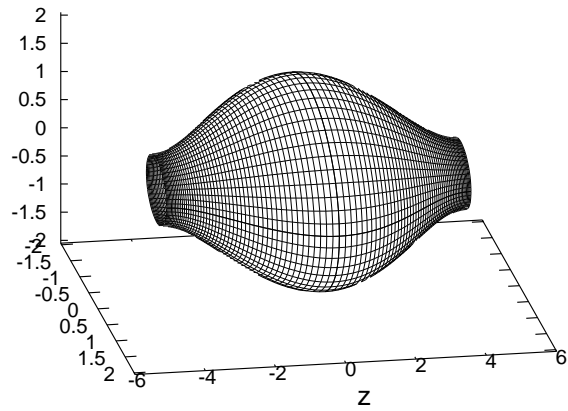


Embedded Horizon

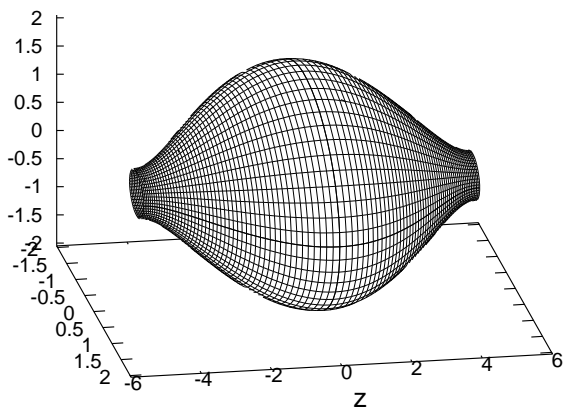
$\lambda=0.5$



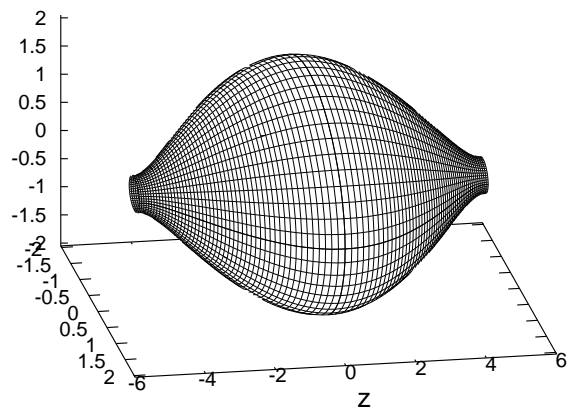
$\lambda=1.0$



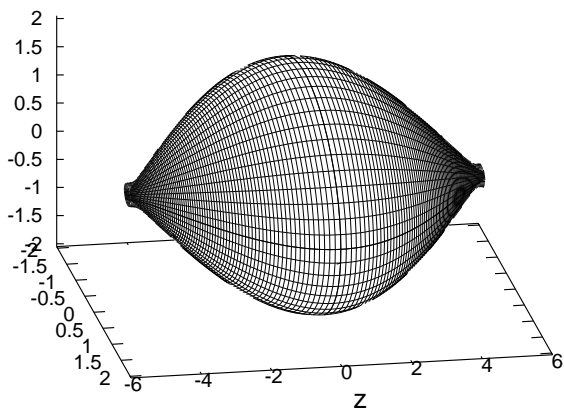
$\lambda=2.0$



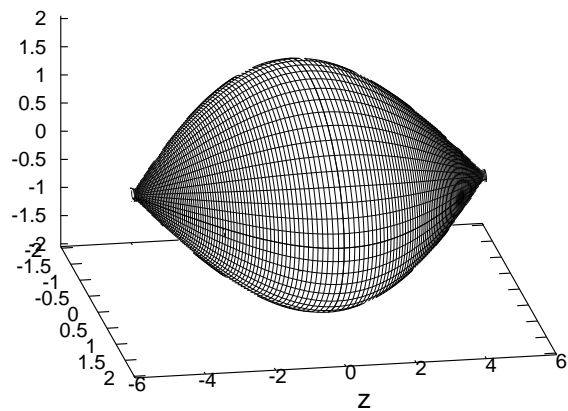
$\lambda=3.0$



$\lambda=5.0$



$\lambda=9.0$



Physical Quantities

- Mass M , Tension \mathcal{T} , relative tension $n = \mathcal{T}/M$
- Temperature $T = e^{d_0} T_0$
- Entropy $S = \frac{S_0}{L} \int_0^L e^{B_H + (D-3)C_H} dz$

$$\text{Smarr relation } TS = \frac{D-3-n}{D-2} M$$

$$1^{\text{st}} \text{ law } dM = TdS + \mathcal{T}dL$$

- Asymptotics ($D = 5$)

$$g_{tt} \rightarrow -1 + \frac{c_t}{r}, \quad g_{zz} \rightarrow 1 + \frac{c_z}{r}$$

$$M = \frac{\Omega_2 L}{16\pi G} (2c_t + c_z), \quad n = \frac{c_t - 2c_z}{2c_t - c_z}$$

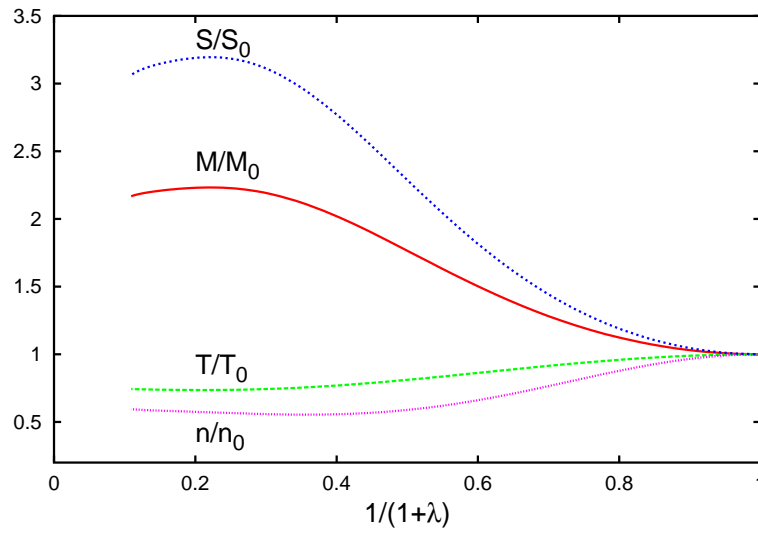
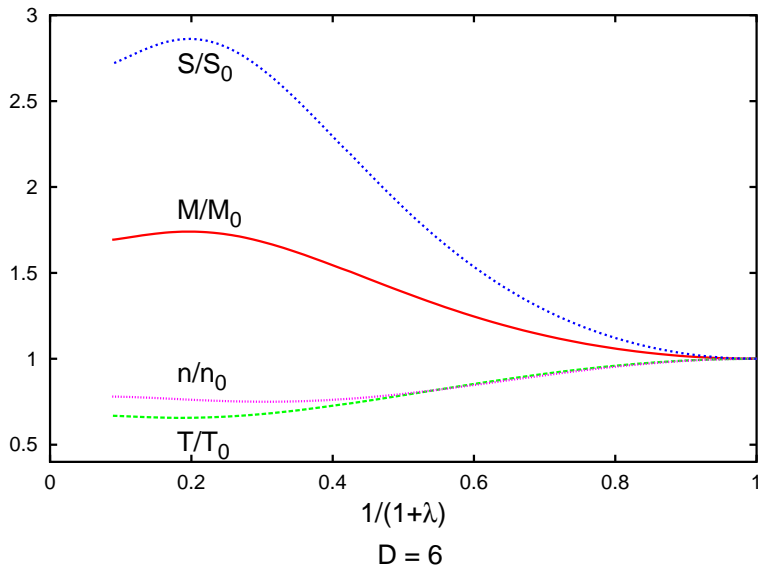
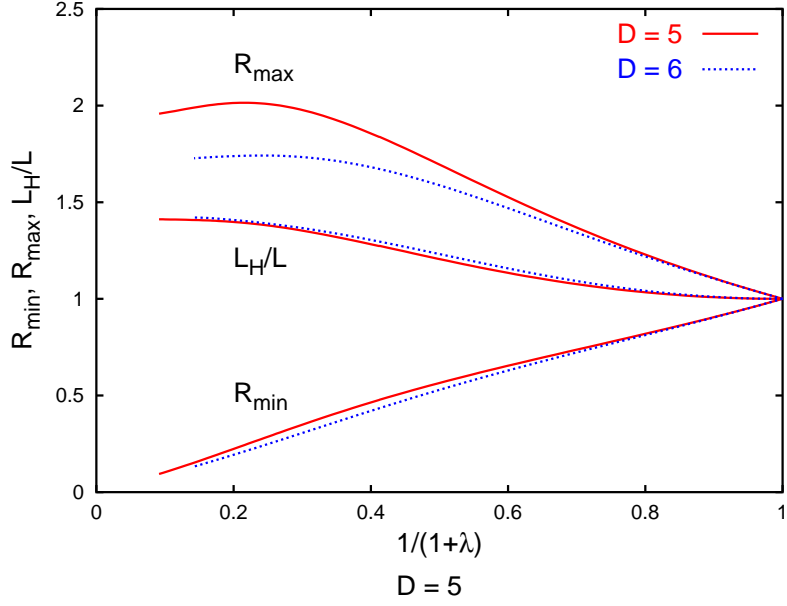
- c_z is error prone,

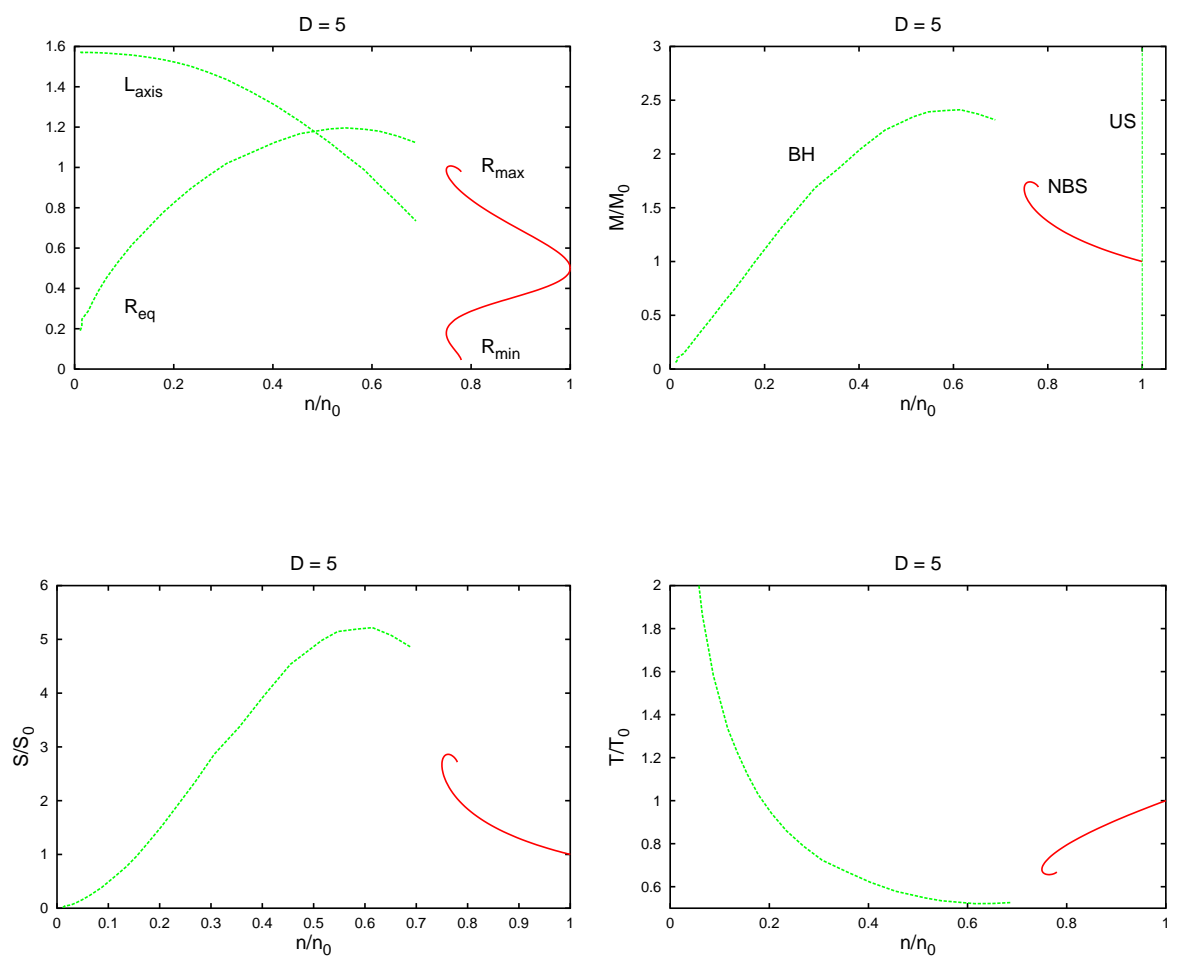
$$\text{use instead } 1^{\text{st}} \text{ law: } M = M_0 + \int_{S_0}^S T(s') ds'$$

and Smarr relation to obtain M, n

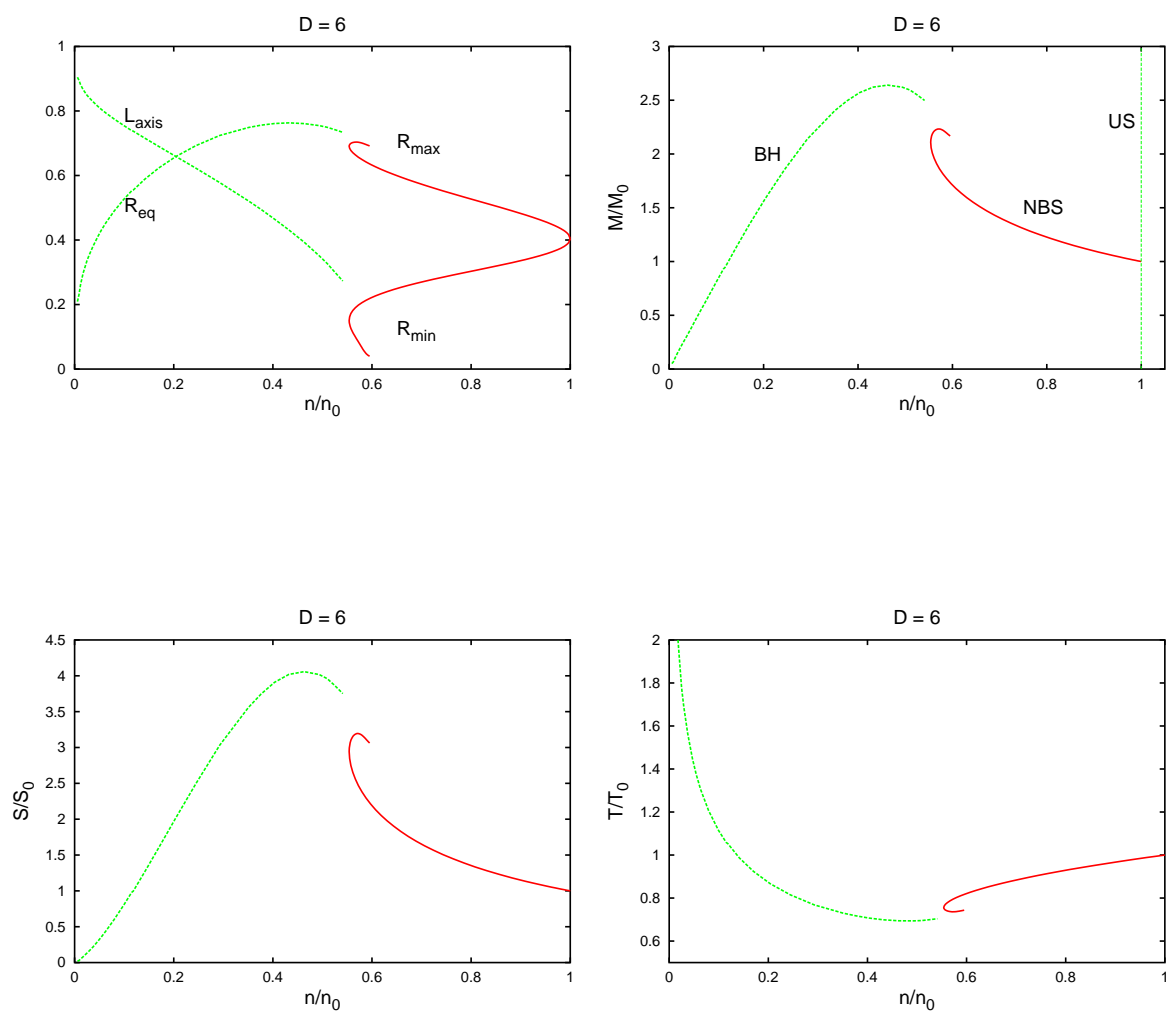
- R_{\max} (R_{\min}) max (min) radius of S^{D-3} on the horizon

L_H proper length of the horizon
in compact direction





black hole data from Kudoh and Wiseman, PRL94(2005)161102



black hole data from Kudoh and Wiseman, PRL94(2005)161102

- **Action**

$$I = \frac{1}{16\pi G} \int d^D x \sqrt{-g} \left(R - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{4} e^{-2a\phi} F^2 \right)$$

$F = dA$, ϕ dilaton, a dilaton coupling

- **Field equations**

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \frac{1}{2} T_{\mu\nu},$$

$$\nabla^2 \phi = -\frac{a}{2} e^{-2a\phi} F^2,$$

$$\partial_\mu (\sqrt{-g} e^{-2a\phi} F^{\mu\nu}) = 0.$$

- **Vacuum solution**

$$ds^2 = -V(x) dt^2 + h_{ij}(x) dx^i dx^j$$

implies family of **EMD solutions**

$$ds^2 = -V(\cosh^2 \beta - \sinh^2 \beta V)^{-2\alpha(D-3)} dt^2 + (\cosh^2 \beta - \sinh^2 \beta V)^{2\alpha} h_{ij} dx^i dx^j,$$

$$A_\mu = \sqrt{2(D-2)\alpha} \frac{\tanh \beta e^{a\phi_0} V}{\cosh^2 \beta - \sinh^2 \beta V} \delta_{\mu t},$$

$$\phi = \phi_0 - 2a(D-2)\alpha \log(\cosh^2 \beta - \sinh^2 \beta V),$$

$$\alpha = (2a^2(D-2) + D-3)^{-1}, \beta, \phi_0 \text{ constant}$$

- Vacuum solution $(M, \mathcal{T}, n, S, T)$

EMD solutions

$$\bar{M} = M(1 + 2(D - 3 - n)\alpha \sinh^2 \beta),$$

$$\bar{\mathcal{T}} = \mathcal{T},$$

$$\bar{n} = \frac{n}{1 + 2(D - 3 - n)\alpha \sinh^2 \beta}.$$

$$\bar{T} = T(\cosh \beta)^{-2\alpha(D-2)},$$

$$\bar{S} = S(\cosh \beta)^{2\alpha(D-2)},$$

- Electric charge:

$$Q_e = -\frac{M(D - 3 - n)}{(D - 4)} \sqrt{\frac{\alpha}{2(D - 2)}}$$

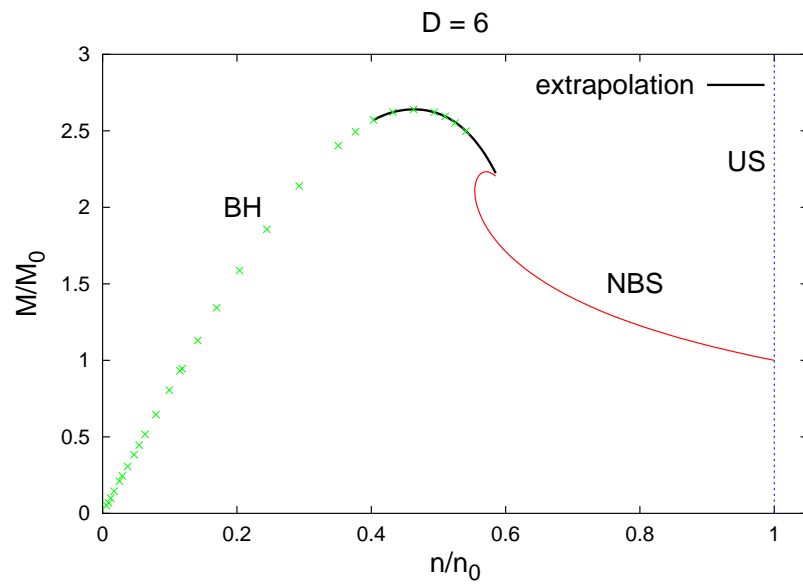
- Dilaton charge:

$$Q_d = -\frac{2aM}{D - 4}(D - 3 - n)\alpha \sinh^2 \beta$$

- Smarr relation:

$$\frac{D - 3 - n}{D - 2}\bar{M} = \bar{T}\bar{S} - \frac{(D - 3)(D - 4)}{D - 2}\Phi\tilde{Q}_e,$$

where $\tilde{Q}_e = \Omega_{D-3}LQ_e$, $\Phi = \sqrt{2(D - 2)\alpha} e^{a\phi_0} \tanh \beta$



black hole data from Kudoh and Wiseman, PRL94(2005)161102

Conclusions

- $D = 5$ and $D = 6$ non-uniform black strings up to $\lambda \approx 9$
- (some) Evidence for topology changing transition
- Black strings in EMD theory from vacuum solutions